

# FINAL EXAM GEOPHYSICAL FLUID DYNAMICS

7 November 2016, 9:00 - 11.00 (2 hours)

Two problems, total 45 points.

Please write your answers for each problem on a separate sheet of paper and put your name and student number on each sheet.

Answers may be written in English or Dutch. Please write clearly and not with a pencil! Remember to write down the correct units of the quantities/numbers you calculate.

In all questions you may use  $g = 10 \text{ ms}^{-2}$ ,  $r_a = 6400 \text{ km}$  and  $\Omega = 7.3 \times 10^{-5} \text{ s}^{-1}$ .

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## Problem 1

Consider the following equation for free surface elevations:

$$\frac{\partial \eta}{\partial t} - R^2 \frac{\partial}{\partial t} \nabla^2 \eta - \beta_0 R^2 \frac{\partial \eta}{\partial x} = 0, \quad (*)$$

- Name and define the parameters  $R$  and  $\beta_0$ .  
Also, explain the physical meaning of these parameters. (5 points)
- Substitute wave-like solutions in Equation (\*) above and show that the dispersion relation of the waves can be recapitulated as

$$\left(k_x + \frac{\beta_0}{2\omega}\right)^2 + k_y^2 = \frac{\beta_0^2}{4\omega^2} - \frac{1}{R^2}.$$

What is the name of these waves? (5 points)

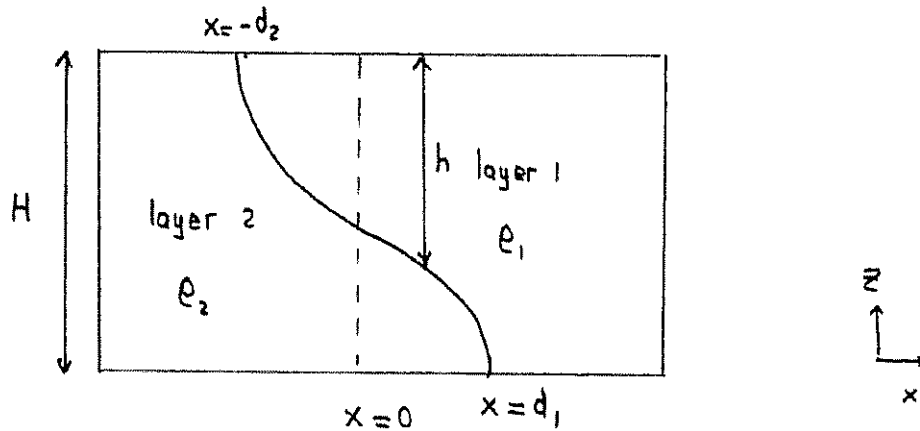
- Consider waves for which  $\omega = \frac{1}{2}\omega_{max}$ , where  $\omega_{max}$  is the maximum frequency that the waves governed by Equation (\*) can attend.  
Derive an expression for  $\omega$  in terms of the parameters  $R$  and  $\beta_0$ . (5 points)
- The waves of item b ( $\omega = \frac{1}{2}\omega_{max}$ ) also have a meridional wavenumber  $k_y = \sqrt{2}/R$  and their energy propagation speed has an eastward component.  
Compute the zonal wavenumber  $k_x$  of these waves. You will find two solutions for  $k_x$ , please argue which one is the one we are looking for here. (5 points)
- Sketch, at a fixed time, the crest and trough lines of the waves of item c in the  $x-y$  plane, Northern Hemisphere.  
Also indicate, in the same figure, with arrows the velocity experienced by fluid particles.  
Explain your answer. (5 points)

**For problem 2: P.T.O.**

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## Problem 2

Two layers of fluids, with densities  $\rho_1$  and  $\rho_2$ , are initially separated by a vertical wall at  $x = 0$  (the dashed line in the figure below). In the  $y$ -direction you can assume uniform conditions.



At time  $t = 0$  the vertical wall is removed and the system adjusts to a final steady state. In the figure, the interface between the two fluids in the end state is indicated by the solid curve (from  $x = -d_2$  to  $x = d_1$ ). Note that thickness  $h = H - a$ . The dynamics is governed by the nonlinear, frictionless shallow water equations for a two-layer system on the  $f$ -plane.

- a. Considering the situation sketch above, is density  $\rho_1$  larger or smaller than  $\rho_2$ ? Motivate your answer. (5 points)

- b. By analysing the equations of motion, it follows that

$$\frac{dv_1}{dx} = f \left[ \frac{h}{H} - 1 \right], \quad \frac{dv_2}{dx} = -f \frac{h}{H},$$

where  $v_1, v_2$  are the velocities in the layers with densities  $\rho_1, \rho_2$ .

Derive this result and explain the principle that is crucial in this respect. (5 points)

- c. A third equation relating  $v_1, v_2$  and  $h$  in the end state is

$$f(v_1 - v_2) = g' \frac{dh}{dx}.$$

Name this balance, and derive it from the given equations of motion (explain your steps). (5 points)

- d. From the equations of items b and c the thickness  $h(x)$  of the end state can be found, using the boundary conditions  $h(x = -d_2) = 0$  and  $h(x = d_1) = H$ . In order to determine the locations  $d_1$  and  $d_2$  two additional constraints are needed.

Describe these constraints, both physically and in terms of mathematical expressions. (5 points)

**END**

# GFD 2016 Equation sheet

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## Continuity and momentum equations: molecular viscous fluid

$$\frac{\partial \rho}{\partial t} + \frac{\partial}{\partial x}(\rho u) + \frac{\partial}{\partial y}(\rho v) + \frac{\partial}{\partial z}(\rho w) = 0$$

$$\begin{aligned}\rho \left( \frac{du}{dt} + f_* w - f v \right) &= - \frac{\partial p}{\partial x} && + \mu \left( \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} + \frac{\partial^2 u}{\partial z^2} \right) \\ \rho \left( \frac{dv}{dt} + f u \right) &= - \frac{\partial p}{\partial y} && + \mu \left( \frac{\partial^2 v}{\partial x^2} + \frac{\partial^2 v}{\partial y^2} + \frac{\partial^2 v}{\partial z^2} \right) \\ \rho \left( \frac{dw}{dt} - f_* u \right) &= - \frac{\partial p}{\partial z} - \rho g && + \mu \left( \frac{\partial^2 w}{\partial x^2} + \frac{\partial^2 w}{\partial y^2} + \frac{\partial^2 w}{\partial z^2} \right)\end{aligned}$$

## Relative circulation and relative vorticity

$$\Gamma = \oint_C \mathbf{u} \cdot d\mathbf{r} = \int_S (\nabla \times \mathbf{u}) \cdot \mathbf{n} dS$$

where  $S$  is the surface enclosed by contour  $C$ .

## Shallow water equations

$$\begin{aligned}\frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} - f v &= -g \frac{\partial \eta}{\partial x} \\ \frac{\partial v}{\partial t} + u \frac{\partial v}{\partial x} + v \frac{\partial v}{\partial y} + f u &= -g \frac{\partial \eta}{\partial y} \\ \frac{\partial \eta}{\partial t} + \frac{\partial}{\partial x}(hu) + \frac{\partial}{\partial y}(hv) &= 0\end{aligned}$$

## Ekman pump (Northern Hemisphere)

$$\bar{w} = \frac{d}{2} \bar{\zeta}, \quad d = \left( \frac{2\nu_E}{f} \right)^{1/2} \quad \text{and} \quad w_{Ek} = \frac{1}{\rho_0 f} \left[ \frac{\partial \tau^y}{\partial x} - \frac{\partial \tau^x}{\partial y} \right]$$

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## Linear shallow water equations for 2-layer model

$$\begin{aligned}
 \frac{\partial u_1}{\partial t} - f v_1 &= -g \frac{\partial \eta}{\partial x}, & \frac{\partial v_1}{\partial t} + f u_1 &= -g \frac{\partial \eta}{\partial y}, \\
 \frac{\partial u_2}{\partial t} - f v_2 &= -g \frac{\partial \eta}{\partial x} - g' \frac{\partial a}{\partial x}, & \frac{\partial v_2}{\partial t} + f u_2 &= -g \frac{\partial \eta}{\partial y} - g' \frac{\partial a}{\partial y}, \\
 \frac{\partial}{\partial t}(\eta - a) + \frac{\partial}{\partial x}(H_1 u_1) + \frac{\partial}{\partial y}(H_1 v_1) &= 0, \\
 \frac{\partial a}{\partial t} + \frac{\partial}{\partial x}((H_2 - b)u_2) + \frac{\partial}{\partial y}((H_2 - b)v_2) &= 0.
 \end{aligned}$$

Characteristic depth:  $\bar{h} = H_1 H_2 / (H_1 + H_2)$ .

## Generalised equation for barotropic planetary/topographic waves

$$\frac{\partial \eta}{\partial t} - R^2 \frac{\partial}{\partial t} \nabla^2 \eta - \beta_0 R^2 \frac{\partial \eta}{\partial x} + \frac{\alpha_0 g}{f_0} \frac{\partial \eta}{\partial x} = 0.$$

## Complex variable

$$\phi \equiv |\phi| e^{i\theta} \equiv \phi_r + i\phi_i, \quad \text{where } |\phi|^2 = \phi_r^2 + \phi_i^2, \quad \tan \theta = \phi_i / \phi_r.$$

## QG Theory for a continuously stratified fluid

$$\begin{aligned}
 \frac{\partial w}{\partial z} &= \frac{1}{\rho_0 f_0^2} \left[ \frac{\partial}{\partial t} \nabla^2 p' + \frac{1}{\rho_0 f_0} J(p', \nabla^2 p') + \beta_0 \frac{\partial p'}{\partial x} \right], \\
 \frac{\partial \rho'}{\partial t} + \frac{1}{\rho_0 f_0} J(p', \rho') - \frac{\rho_0 N^2}{g} w &= 0, & N^2 &= -\frac{g}{\rho_0} \frac{d\bar{\rho}}{dz}, \\
 \frac{\partial q}{\partial t} + J(\psi, q) &= 0, & q &= \nabla^2 \psi + \frac{\partial}{\partial z} \left( \frac{f_0^2}{N^2} \frac{\partial \psi}{\partial z} \right) + \beta_0 y.
 \end{aligned}$$

## Two-layer QG model

$$\begin{aligned}
 q_1 &= \nabla^2 \psi_1 + \frac{1}{2R^2} (\psi_2 - \psi_1) + f_0 + \beta_0 y, \\
 q_2 &= \nabla^2 \psi_2 - \frac{1}{2R^2} (\psi_2 - \psi_1) + f_0 + \beta_0 y, & R &= \frac{\sqrt{g'H}}{2f_0}, \\
 w &= \frac{2f_0}{N^2 H} \left[ \frac{\partial}{\partial t} (\psi_2 - \psi_1) + J(\psi_1, \psi_2) \right], & N^2 &= \frac{2g'}{H}, \\
 a &= \frac{f_0}{g} (\psi_2 - \psi_1), & \rho' &= \rho_0 \frac{f_0}{2gH} (\psi_2 - \psi_1).
 \end{aligned}$$


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