

**Final Exam Advanced Quantum Mechanics (total 300 points)**

Tuesday, January 30, 2018, 13:30-16:30

1. Write your name and initials on all sheets, on the first sheet also your student ID number.
2. Write clearly, unreadable work cannot be corrected.
3. Give the motivation, explanation, and calculations leading up to each answer and/or solution.
4. Do not spend a large amount of time on finding (small) calculational errors. If you suspect you have made such an error, point it out in words.
5. Note the appendix at the end of this exercise!
6. This exam consists of three exercises. Start each exercise on a new sheet of paper.

**1. SOME CONCEPTS (80 POINTS)**

- a) (30 points) Describe in a few lines — using a few equations if you want — what is meant by i) entanglement, ii) a mixed state, iii) a pure state, and iv) decoherence.
- b) (10 points) Describe what a superposition is and/or give an example of a superposition.
- c) (20 points) Describe briefly what a measurement of violation of the Bell inequalities implies (NB: do not describe how such measurements are done, but, essentially, why they are done).
- d) (10 points) Give an example of a practical application of entanglement that is currently envisioned, or is already in use.
- e) (10 points) Consider a spin with total spin angular momentum quantum number  $S_1 = 4$ , and another spin with total spin angular momentum quantum number  $S_2 = 3/2$ . Give the possible outcomes of a measurement of the total spin angular momentum of a quantum system that consists of both these spins together.

**2. PARTICLE WITH SPIN IN A ROTATING FRAME (100 POINTS)**

Consider a two-dimensional system described by coordinates  $(x, y)$ , consisting of one spin-1/2 particle ( $S = 1/2$ ) with mass  $m$  in a potential  $V(x, y)$  that depends only on  $r = \sqrt{x^2 + y^2}$ . In the basis-independent formulation, the time-dependent Schrödinger equation for the state vector  $|\Psi\rangle$  is therefore

$$i\hbar \frac{\partial}{\partial t} |\Psi\rangle = \left[ \frac{\hat{p}_x^2 + \hat{p}_y^2}{2m} + V(\hat{x}, \hat{y}) \right] |\Psi\rangle .$$

Consider now a transformation to a coordinate frame rotating, with angular frequency  $\omega$  around the  $z$ -axis, with respect to the original frame.

a) (20 points) Show that the Schrödinger equation in this rotating frame is

$$i\hbar \frac{\partial}{\partial t} |\Psi\rangle = \left[ \frac{\hat{p}_x^2 + \hat{p}_y^2}{2m} - \omega (\hat{L}_z + \hat{S}_z) + V(\hat{x}, \hat{y}) \right] |\Psi\rangle,$$

with  $\hat{L}_z$  the operator corresponding to the orbital angular momentum in the  $z$ -direction, and  $\hat{S}_z$  the  $z$ -component of the spin angular momentum.

b) (30 points) Define the wave functions  $\Psi_\sigma(x, y) \equiv (\langle x, y | \otimes \langle \sigma |) |\Psi\rangle$ , with  $|x, y\rangle$  the eigenstates of the position operator and  $|\sigma\rangle$  the eigenstates of  $\hat{S}_z$  with respective eigenvalues  $\sigma\hbar/2$ , and  $\sigma \in \{+1, -1\}$ . Show that the Schrödinger equation in position representation is given by

$$i\hbar \frac{\partial \Psi_\sigma(x, y, t)}{\partial t} = \left[ -\frac{\hbar^2}{2m} \left( \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} \right) - \frac{\sigma\hbar\omega}{2} + i\hbar\omega \left( x \frac{\partial}{\partial y} - y \frac{\partial}{\partial x} \right) + V(x, y) \right] \Psi_\sigma(x, y, t).$$

c) (30 points) Show that this latter Schrödinger equation can be written in the form of that corresponding to a particle with charge  $q$  in an electromagnetic field so that

$$i\hbar \frac{\partial \Psi_\sigma(x, y, t)}{\partial t} = \left[ \frac{(\hat{\mathbf{p}}_{\text{pos}} - q\mathbf{A})^2}{2m} + V(x, y) - \frac{\sigma\hbar\omega}{2} + V_\omega(x, y) \right] \Psi_\sigma(x, y, t),$$

with  $\hat{\mathbf{p}}_{\text{pos}} = -i\hbar(\partial/\partial x, \partial/\partial y, 0)$  the momentum operator in the position representation,  $\mathbf{A} = m\omega\mathbf{e}_z \times (x, y, 0)/q$  an effective vector potential with  $\mathbf{e}_z$  the unit vector in the  $z$ -direction, and  $V_\omega(x, y) = -m\omega^2(x^2 + y^2)/2$ .

d) (20 points) The above shows that the dynamics of a particle in a rotating frame is effectively the same as that of a charged particle in a magnetic field  $\mathbf{B} = \nabla \times \mathbf{A}$ . i) Give a physical reason why this is so, and give the expression for  $\mathbf{B}$ . ii) Give a physical interpretation of the potential  $V_\omega(x, y)$ .

### 3. ATOMIC NUCLEUS IN AN ELECTRIC FIELD (120 POINTS)

Consider a particle with spin  $s$ . We ignore its orbital motion and assume its normalized state vector to be

$$|\Psi\rangle = \sum_{m_s=-s}^s a_{m_s} |s, m_s\rangle$$

Here, the states  $|s, m_s\rangle$  are the simultaneous eigenstates of  $\hat{S}_z$  and  $\hat{S}^2$  and the coefficients  $a_{m_s}$  are complex numbers.

a) (15 points) For the case  $s = 1$ , what are the possible outcomes for a measurement along the  $y$ -direction and what are the corresponding probabilities? (NB: Note the appendix at the end of the exam!)

Assume now that an ensemble of such spin  $s$  particles is described by the density operator

$$\hat{\rho} = \sum_{m_s=-s}^s p_{m_s} |s, m_s\rangle \langle s, m_s|$$

b) (15 points) Give the physical meaning of the numbers  $p_{m_s}$  and at least one condition which these numbers must satisfy for the case of general  $s$ . Give the density operator for the case  $s = 1$  for i) a pure state, and ii) a uniformly distributed ensemble, and describe their physical meanings.

Consider a Hamiltonian of the form [1]

$$\hat{H} = k \sum_{i,j=1}^3 \hat{S}_i V_{ij} \hat{S}_j$$

where  $k$  is a constant,  $\hat{S}_i$  is the spin operator corresponding to the  $i$ -th Cartesian direction and  $V_{ij}$  are the elements of a real traceless symmetric  $3 \times 3$  matrix  $V$  — note that this latter matrix is not a quantummechanical operator.

c) (15 points) Why can you always rotate to new coordinates  $x, y, z$  in which  $V$  is diagonal?

Hence, in those new coordinates,

$$V = \begin{pmatrix} V_{xx} & 0 & 0 \\ 0 & V_{yy} & 0 \\ 0 & 0 & V_{zz} \end{pmatrix}$$

with  $V_{xx} + V_{yy} + V_{zz} = 0$ .

d) (20 points) Show that the Hamilton operator in these coordinates can be written as

$$\hat{H} = K \left( 3\hat{S}_z^2 - \hat{S}^2 + \eta(\hat{S}_x^2 - \hat{S}_y^2) \right), \quad \eta = \frac{V_{xx} - V_{yy}}{V_{zz}},$$

and give  $K$ . Here, we assumed that  $V_{zz} \neq 0$ .

e) (10 points) Give the exact eigenvalues and eigenvectors of  $\hat{H}$  for the case  $\eta = 0$  and any spin  $s$ . How many energy levels are there and what is the degree of degeneracy of each level?

f) (30 points) Find the exact eigenvalues and normalized eigenvectors of  $\hat{H}$  for general  $\eta$  and spin 1 and draw a graph of the energy levels as function of  $\eta$  for  $0 \leq \eta \leq 1$ . Hint: Start by expressing  $\hat{S}_x^2 - \hat{S}_y^2$  in terms of  $\hat{S}_\pm = \hat{S}_x \pm i\hat{S}_y$ .

g) (15 points) Consider such a spin 1 in thermal equilibrium at temperature  $T$ . Find the expectation value of  $\hat{S}_z$  for  $\eta = 0$ .

Appendix — Possibly, you may want to make use of these results:

- Under a rotation of angle  $\theta$  around an axis along the unit vector  $\mathbf{n}$ , the state vector transforms according to  $|\Psi\rangle \rightarrow e^{-i\theta \mathbf{n} \cdot \hat{\mathbf{J}}} |\Psi\rangle$ , where  $\hat{\mathbf{J}}$  is the operator corresponding to total angular momentum.
- For an angular momentum operator  $\hat{\mathbf{J}}$  we have the states  $|j, m\rangle$ , which are the eigenstates of  $\hat{\mathbf{J}}^2$  and  $\hat{J}_z$ . The respective eigenvalues are  $j(j+1)\hbar^2$  and  $m\hbar$ , where  $m$  runs from  $-j$  to  $+j$  in integer steps and  $j$  is an integer. Furthermore, we have that  $\hat{J}_\pm = \hat{J}_x \pm i\hat{J}_y$  which act on these states as:

$$\hat{J}_+ |j, m\rangle = \hbar \sqrt{j(j+1) - m(m+1)} |j, m+1\rangle;$$

$$\hat{J}_- |j, m\rangle = \hbar \sqrt{j(j+1) - m(m-1)} |j, m-1\rangle.$$

- For  $s = 1$  and in the basis  $\{|1, 1\rangle, |1, 0\rangle, |1, -1\rangle\}$ , the spin operators  $\hat{S}_i$  are represented by

$$\hat{S}_x = \frac{\hbar}{\sqrt{2}} \begin{pmatrix} 0 & 1 & 0 \\ 1 & 0 & 1 \\ 0 & 1 & 0 \end{pmatrix}, \quad \hat{S}_y = \frac{\hbar}{\sqrt{2}} \begin{pmatrix} 0 & -i & 0 \\ i & 0 & -i \\ 0 & i & 0 \end{pmatrix}, \quad \hat{S}_z = \hbar \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & -1 \end{pmatrix}$$

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[1] For example, an atomic nucleus with a non-zero electric quadrupole moment in an inhomogeneous electric field is described by this hamiltonian. In that case  $V$  results from the gradient of the electric field.