
Retake Exam for Advanced Statistical Physics: NS-370B

Date: December 18th, 2023

Time for Regular Students: 18:30 - 21:30

Extra-Time Students Only: 18:30 – 22:00

This exam consists of **6** questions of varying length.

The total number of possible points is: 100.

This is a **closed-book** exam, no notes are allowed.

You will not require a calculator or any other electronic device and are consequently not permitted one.

Please **start every exercise on a new PAGE of paper, NOT a new 4-PAGE SHEET.**

Write your name clearly on every four-page sheet of paper.

Please number your sheets (1/3, 2/3, ...) and indicate which questions are on it.

The exam can be written in English or Dutch.

Please write CLEARLY and using PEN!

Equation Cheat Sheet

A few formulas and information that may be useful for this exam:

- The canonical partition function of a classical thermodynamic system of N identical particles in a volume V a temperature T with Hamiltonian $H(\mathbf{\Gamma})$ — $\mathbf{\Gamma}$ is a point in phase space — is written:

$$Z(N, V, T) = \frac{1}{N!h^{3N}} \int d\mathbf{\Gamma} \exp[-\beta H(\mathbf{\Gamma})], \text{ where } \beta^{-1} = k_B T.$$

- The grand partition function of identical particles is given by (pay attention to the N dependence of Z for each term in the sum):

$$\Xi(\mu, V, T) = \sum_{N=0}^{\infty} \exp(\beta\mu N) Z(N, V, T),$$

and the grand-canonical distribution is given by:

$$f_g(\mathbf{\Gamma}, N) = \frac{\exp[\beta\mu N - \beta H(\mathbf{\Gamma})]}{N!h^{3N}\Xi(\mu, V, T)}, \text{ with } \mu \text{ the chemical potential.}$$

- The second virial coefficient B_2 for an isotropic pair potential $\phi(r)$ is given by:

$$B_2 = -\frac{1}{2} \int_V d\mathbf{r} (\exp(-\beta\phi(r)) - 1).$$

- $k_B = 1.38 \times 10^{-23}$ J/K, $e = 1.6 \times 10^{-19}$ C, and $R = 8.31$ J/K/mol.

- The binomial coefficient (*i.e.*, m choose n), which is the number of ways n objects can be chosen from m objects, is given by:

$$\binom{m}{n} = \frac{m!}{(m-n)!n!}.$$

- Stirling's approximation to order $O(N)$ is given by: $\log(N!) = N \log N - N$.

- Gaussian Integral:

$$\int_{-\infty}^{\infty} dx e^{-x^2} = \sqrt{\pi}.$$

- The Taylor series of $f(x)$ around $x = a$ is given by:

$$f(x) = f(a) + \frac{f'(a)}{1!}(x-a) + \frac{f''(a)}{2!}(x-a)^2 + \frac{f^{(3)}(a)}{3!}(x-a)^3 + \dots$$

- From the Taylor series we obtain:

$$e^x = \sum_{n=0}^{\infty} \frac{x^n}{n!} \quad \& \quad \log(1-x) = -\sum_{n=1}^{\infty} \frac{x^n}{n}.$$

- Maxwell-Boltzmann ($s = 0$), Bose-Einstein ($s = -1$), and Fermi-Dirac ($s = 1$) statistics are:

$$f(\epsilon) = \frac{1}{\exp[\beta(\epsilon - \mu)] + s},$$

where ϵ is the energy (of a single-particle level) and μ the chemical potential.

- The quantum single-level 'grand'-canonical partition functions Z_i are given by:

$$Z_i = \sum_{n_i=0}^M \exp(\beta n_i(\mu - \epsilon_i)),$$

where ϵ_i is the energy associated with the i -th level and μ is the chemical potential. The value of M depends on the statistics; $M = 1$ for fermions and $M = \infty$ for bosons. These lead to the following grand canonical partition function

$$\Xi(\mu, V, T) = \prod_i Z_i,$$

for the system, where i runs over all available levels.

1. Thermodynamics: (16 points)

In this question, we consider several useful thermodynamic relations, identities, and transformations.

- Write down the mathematical expressions for the *first* and *second* law of thermodynamics. Write down what the symbols mean and explain the sign convention that you picked in a *few* words.
- Provide the two standard forms of work. Ensure agreement with your choice in part (a).
- Combine the laws and choices of work to obtain dU . Explain in a *few* words what must hold to use this expression.
- Make the transition to the μVT ensemble by deriving $d\Omega = -SdT - pdV - Nd\mu$. Provide the transformation **and** the intermediate steps.
- Derive expressions for N , p , and S in terms of $\Omega(\mu, V, T)$ based on your result from (d).
- Show that $\Omega(\mu, V, T) = -p(\mu, T)V$ and explain in *few* words.

2. Gibbs and Boltzmann Entropy: (16 points)

The Gibbs definition of entropy S states that $S = -k_B \sum_s p_s \log p_s$, where the sum is over all states labelled by s . The probability of being in that state is p_s and k_B is Boltzmann's constant.

- Show that the Gibbs definition leads to an extensive entropy.

Next, we will use $\sum_s p_s = 1$ and $\langle \epsilon_s \rangle = E$ as a constraint to maximize S .

- Define the Lagrangian \mathcal{L} and show that $p_s \propto \exp(-c\epsilon_s)$, with c some constant.
- Use the definition of temperature $1/T = (\partial S/\partial U)_{N,V}$ to determine c . Under what conditions do we set $U = E$ typically? Explain in a *few* words.

3. Small Clusters: (21 points)

Consider a square-lattice model with A sites of which n are occupied by particles. If two particles are adjacent, they bond with bonding energy $-\epsilon$ (with $\epsilon > 0$). However, a particle that is already bonded cannot bond with any other, *i.e.*, the system consists of monomers and dimers. Here, we develop an approximate solution to the problem for low densities. Further assume that there is an average monomer density $\rho_1 = n_1/A$ and dimer density $\rho_2 = n_2/A$, with n_1 and n_2 the number of monomers and dimers, respectively. Also assume that the monomers and dimers do not interact, so both behave ideal gas like.

- Show that the canonical partition sum for the monomers $Q_1(n_1, A, T) = A^{n_1}/(n_1!)$, with T the temperature. Assume periodic boundary conditions and explain your answer using a *few* words.
- Use your answer in (a) to provide the canonical partition function for the dimers. Hint: Remember that the dimers can have two orientations.
- Calculate the (reduced) Helmholtz free energy

$$\beta F(\rho_1, \rho_2, A, T)/A = \rho_1(\log \rho_1 - 1) + \rho_2(\log \rho_2 - 1) - \rho_2(\log 2 + \beta\epsilon), \quad (1)$$

of the system, where $\beta = 1/(k_B T)$ is the inverse thermal energy with k_B Boltzmann's constant. Hint: Use Stirling's approximation $\log n! = n \log n - n$.

- Show that in equilibrium the chemical potentials of monomers and dimers should obey: $2\mu_1 = \mu_2$.
- Demonstrate that $2\mu_1 = \mu_2$ imposes $\rho_2 = 2\rho_1^2 \exp(\beta\epsilon)$.
- Use the relation between n , n_1 , and n_2 to arrive at

$$\rho_1 = \frac{\sqrt{1 + 16\rho_2 e^{\beta\epsilon}} - 1}{8e^{\beta\epsilon}}. \quad (2)$$

- What happens to ρ_1 and ρ_2 in the limit where $\epsilon \rightarrow 0$ (assuming low density)? Hint: You can either perform the lengthy calculation or use your physical intuition. In either case explain your answer in a *few* words.

The exam continues on the next page \Rightarrow

4. Percolation: (7 points)

Consider a square lattice. Assume that a 2×2 subset in this lattice ‘spans’ if there is a vertically spanning cluster. Show that $R(p) = 2p^2(1-p)^2 + 4p^3(1-p) + p^4$ and support your argumentation toward this result using sketches. Determine the fully reduced analytic expression for the non-trivial fixed point.

5. Bose Gas at High Temperatures: (23 points)

At high temperatures T , a gas of conserved bosons in a volume V behaves similar to a classical ideal gas. In this problem, the goal is to determine the leading-order corrections to the classical ideal gas law $pV = Nk_B T$, where p is the pressure, N the particle number, and k_B is the Boltzmann constant. We start from the grand-canonical partition function $\Xi(\mu, V, T) = \prod_i Z_i$ with $Z_i = \sum_{n_i=0}^{\infty} \exp(\beta n_i(\mu - \epsilon_i))$, where μ is the chemical potential and ϵ_i the energy of the i -th quantum level of a single particle.

- (a) Use a general relation between $\langle N \rangle$ and Ξ to show that the particle number N is given by the integral representation

$$N = \int dE g(E) f_{\text{BE}}(E), \quad (3)$$

where $g(E)$ is the density of states and $f_{\text{BE}}(E)$ adheres to Bose-Einstein statistics.

- (b) Use the definition of the wave vector $\mathbf{k} = \pi \mathbf{n} / L$, with \mathbf{n} the vector containing the quantum numbers n_x, n_y , and n_z , and the expression for the energy $E_{\mathbf{n}} = (\hbar^2 k^2) / (2m)$, with m the particle mass, $k = |\mathbf{k}|$, and $\hbar = h / (2\pi)$ the reduced Planck constant, to arrive at the density of states

$$g(E) = \frac{L^3}{4\pi^2} \left(\frac{2m}{\hbar^2} \right)^{3/2} \sqrt{E}. \quad (4)$$

We will identify $L^3 = V$ for the remainder of this exercise.

- (c) Input both $g(E)$ and $f_{\text{BE}}(E)$ into the expression for N and rework it using a substitution to arrive at the following scaling of the particle density $\rho = N/V$ with the thermal wavelength Λ

$$\rho = \frac{1}{\Lambda^3} \frac{2}{\sqrt{\pi}} \int dx \frac{\sqrt{x}}{z^{-1}e^x - 1}, \quad (5)$$

with $z = \exp(\beta\mu)$. Hint: You can find the expression for $f_{\text{BE}}(E)$ on the cheat sheet if you could not solve part (a).

- (d) Under what criterion involving Λ and ρ will the system show significant quantum effects? Explain using a *few* words.
- (e) Argue using only a *few* words and the above expression for N that at high temperatures T the chemical potential μ becomes large but negative, such that $z = e^{\beta\mu} \ll 1$. In your answer you should reference the ensemble that you must choose to work in to make the comparison to the classical ideal-gas law.
- (f) Turning to a **classical** ideal gas (in 3D), use equipartition to show that the pressure satisfies $pV = 2E/3$ with E the energy. Explain your answer in a *few* words.

Using your answer in (c) and an analogous expression for E it can be shown that the quantum correction to the classical ideal gas result is given to leading-order by

$$\beta p = \rho - \frac{\Lambda^3}{4\sqrt{2}} \rho^2. \quad (6)$$

- (g) Does this result make sense? Explain in a *few* words whether you would expect to have a lower pressure for bosons than purely classical particles at the same density.

The exam continues on the next page \Rightarrow

6. Landau Theory for Needle-like Objects: (17 points)

Here, we examine the phase behavior of anisotropic particles (long rods) using Landau theory for the isotropic-nematic (IN) phase transition without an external field.

- (a) What are the isotropic and nematic phase for a (colloidal) liquid crystal? Illustrate using sketches and explain using a *few* words only.
- (b) Provide an approximate expression for the second virial coefficient B_2 for thin, long rods of length L and diameter D , *i.e.*, $D \ll L$. Explain your expression using a *few* words. Hint: The expression for B_2 can be found on the equation cheat sheet; it should depend on the respective orientation of the rods ($\hat{\omega}_1$ and $\hat{\omega}_2$).

The nematic director is the average orientation of the rods in the liquid crystal. Let θ measure the angle between a rod-like colloid and this director. It is sensible to create an order parameter that depends on a series in terms of $\cos \theta$, which takes values in the range $[-1, 1]$.

- (c) Show that $S = \langle (1/2) (3 \cos^2 \theta - 1) \rangle$ is a suitable order parameter for the disordered phase, using its properties and the requirements of Landau theory.

A physicist arrives at the following form for the IN Landau free energy

$$F_{\text{IN}} = [a(T - T_c) - b^2/(4c)] S^2 + cS^2 [S - b/(2c)]^2, \quad (7)$$

where a , b , and c are positive constants, T_c indicates the temperature for which the isotropic phase fully destabilizes, and T is the temperature of the system. This vanishes $S = 0$ and for $S = b/(2c)$, provided the temperature is suitably chosen.

- (d) What is the order of the isotropic-nematic phase transition? Use arguments supported by the above Landau theory and only a *few* words. You may want to provide a sketch.
- (e) Explain in a *few* words how the hard-rod F may be used to obtain the free energy for a system of rod-like particles with a small, short-ranged attraction. Assume that the attraction can be seen as a perturbation. Provide the expression that connects the source and target free energies. **Hint: No more than one equation is needed!**

The exam ends here ☺!