

$$1 \text{ L} = 1 \text{ dm}^3 \text{ water is } 1 \text{ kg} \quad \frac{1 \text{ kg}}{1 \text{ dm}^3} = \frac{1 \cdot 10^{-3} \text{ kg}}{1000 \frac{\text{m}^3}{\text{m}^3}}$$

Midterm Exam **Soft Condensed Matter Theory**, April 11, 2018, 13:30h-16:30. This exam consists of 18 items, the maximum score for each item is 5 points. Write your name on each page. This is a *closed-book* exam, and electronic tools are **not** allowed. Give arguments for your answers and write clearly -unreadable answers are no answers. You may use that the viscosity of water is $\eta = 10^{-3} \text{ Pa s}$, the Bjerrum length of water at room temperature is 0.72 nm , the Stokes-Einstein equation for the diffusion coefficient of a sphere of radius a reads $D = k_B T / 6\pi\eta a$ with T the temperature and $k_B = 1.38 \cdot 10^{-23} \text{ J/K}$ the Boltzmann constant. The differential of the internal energy is $dU = TdS - pdV + \mu dN + \gamma dA + \psi dQ - fdL + \dots$ with the usual meaning of symbols.

Problem 1

Consider a one-component system in a volume V at temperature T and chemical potential μ . The number of particles in this system, N , fluctuates according to the probability distribution $W(N) = \Xi^{-1}(\mu, V, T) \exp[(\mu N - F(N, V, T))/k_B T]$ where $F(N, V, T)$ is the Helmholtz free energy of the N -particle system and $\Xi(\mu, V, T)$ the grand-canonical partition function.

- ⌋ (a) Show that the average number of particles, $\langle N \rangle$, is equal to $k_B T (\partial \ln \Xi(\mu, V, T) / \partial \mu)_{V, T}$.
- (b) Show first that $k_B T (\partial \langle N \rangle / \partial \mu)_{V, T} = \langle N^2 \rangle - \langle N \rangle^2$, and use then e.g. the Gibbs-Duhem equation to derive that $k_B T (\partial \rho / \partial p)_T = (\langle N^2 \rangle - \langle N \rangle^2) / \langle N \rangle$, with p the pressure.

Consider now the case of a homogeneous bulk fluid of colloidal hard spheres with a diameter σ at density ρ and temperature T , with the Carnahan-Starling excess (over ideal) free energy $F_{exc} = N k_B T (4\eta - 3\eta^2) / (1 - \eta)^2$ with η the packing fraction.

- ⌋ (c) Give an expression for η in terms of the given variables and calculate the hard-sphere pressure p within the Carnahan-Starling approximation.
- (d) Calculate the second-virial coefficient $B_2(T)$ of hard spheres and check whether or not your result of (c) is consistent with $B_2(T)$.
- ⌋ (e) Sketch the radial distribution function $g(r)$ of these hard spheres for (i) $\eta = 0.01$ and (ii) $\eta = 0.49$. Think of units/scales on both axes and distinguish the two graphs clearly.
- ⌋ (f) Derive and/or explain how the hard-sphere free energy $F_{HS}(N, V, T)$ and the hard-sphere radial distribution function $g(r; \eta)$ can be used to accurately calculate the free energy of dense Lennard-Jones fluids.
- (g) What happens to a colloidal hard-sphere system if $\eta > 0.5$? Briefly describe two experimental techniques with which this can be observed.
- (h) An experimentalist tracks many trajectories of a single colloidal sphere in a dilute hard-sphere suspension, and finds the displacements in the cartesian x -direction after time t to be distributed as a Gaussian $\propto \exp(-\alpha x^2/t)$ with $\alpha = 10 \text{ s}/\mu\text{m}^2$. Estimate the diameter of these hard spheres if the medium is water at room temperature.

Problem 2

The Helmholtz free energy $F(N, V, T)$ of N Argon atoms of mass $m = 6.6 \cdot 10^{-26} \text{ kg}$ in a volume V (density $\rho = N/V$) at temperature T is assumed to be given by

$$F(N, V, T) = N k_B T \left(\log \frac{N \Lambda^3}{V - N b} - 1 \right) - \frac{a N^2}{V}, \quad (1)$$

with constants $a > 0$ and $b > 0$, and Λ the thermal DeBroglie wavelength. Here k_B is the Boltzmann constant.

- (a) Give a physical interpretation of a and b , and give their dimensionality.
- (b) Derive an expression for Λ and explain whether $b \gg \Lambda^3$ or $b \ll \Lambda^3$ at $T = 300\text{K}$.
- (c) Sketch the free energy density F/V as a function of $\rho \in [0, 1/b]$ for (i) $T > T_c \equiv 8a/27bk_B$ and (ii) $T < T_c$, and graphically construct the coexisting gas- and liquid density ρ_g and ρ_l , respectively. Explain your construction briefly.

We now consider an interface between the coexisting bulk gas phase (with density ρ_g) and bulk liquid phase (with density ρ_l). The interface height at horizontal position (x, y) is denoted by the height $z = h(x, y)$, with a gravitational force mg acting on every Argon atom in the downward vertical z -direction. The interfacial tension is known and denoted by γ , and the average height can be set to zero without loss of generality.

- (d) Present a derivation, or give and explain the ingredients, of the capillary wave Hamiltonian $H_{cw} = \frac{\gamma}{2} \int dx dy [(\partial h(x, y)/\partial x)^2 + (\partial h(x, y)/\partial y)^2 + \ell^{-2} h^2(x, y)]$ for this interface. Derive an expression for the capillary length ℓ .
- (e) Estimate (with arguments), or calculate, for the case of a typical air-liquid or oil-water interface at room temperature, the order of magnitude of (i) the tension γ and (ii) the capillary length ℓ . Also give the order of magnitude of (iii) the interface thickness $\sqrt{\langle \bar{h}^2 \rangle}$ with \bar{h}^2 the lateral average of h^2 and $\langle \dots \rangle$ the thermal average.

Problem 3

Consider a 1:1 electrolyte (dielectric constant ϵ , temperature T , ion point charges $\pm e$) in the space between two parallel planar electrodes with equal surface potential ψ_s at separation H . The cartesian coordinate normal to the electrodes is $z \in [-H/2, H/2]$, i.e. $z = 0$ is the symmetry midplane of the system. The electrolyte is in osmotic equilibrium with a bulk 1:1 electrolyte with total ion concentration $2\rho_s$ at zero potential. We wish to calculate the electric potential $\psi(z)$ and the ionic concentration profiles $\rho_{\pm}(z)$ within (linearised) Poisson-Boltzmann theory, for $z \in [-H/2, H/2]$. In SI units the Poisson equation reads $\psi''(z) = -Q_e(z)/(\epsilon_0 \epsilon)$ with ϵ_0 the vacuum permittivity and $Q_e(z)$ the total charge density.

- (a) For $z \in (-H/2, H/2)$, write the ion concentrations as Boltzmann distributions, express $Q_e(z)$ in terms of $\rho_{\pm}(z)$, and combine all this into the Poisson-Boltzmann equation for $\psi(z)$.
- (b) Show for $z \in (-H/2, H/2)$ that a dimensionless, properly scaled and sufficiently small potential satisfies $\phi''(z) = \kappa^2 \phi(z)$. Give the relation between $\psi(z)$ and $\phi(z)$ and derive an expression for κ .
- (c) Solve, using the boundary conditions, the equation of (b) for $\phi(z)$ and calculate the surface charge density of the electrodes.

A mixture of N colloidal spheres at positions \mathbf{R}_i and N_s solvent molecules at positions \mathbf{r}_j has potential energy $\Phi(\{\mathbf{R}\}, \{\mathbf{r}\})$ composed of colloid-colloid, colloid-solvent, and solvent-solvent interactions, and has effective solvent-induced interactions $\Phi_{eff}(\{\mathbf{R}\}; \mu_s, T)$ between the colloidal spheres with μ_s the solvent chemical potential.

- (d) Express $\Phi_{eff}(\{\mathbf{R}\}; \mu_s, T)$, with motivation, in terms of $\Phi(\{\mathbf{R}\}, \{\mathbf{r}\})$.
- (e) Give two examples of medium-induced interactions between colloidal spheres.