

1) a) First law $dU = TdS - pdV + \mu dN$ (given)
 so $dG = d(U - TS + pV) =$
 $= -SdT + Vdp + \mu dN$ (4) has T, p, N as independent variables
 so $G = G(T, p, N)$
 is appropriate

Extensivity: $G(N, p, T) = \mu(p, T) N$
 so $dG = \mu dN + N d\mu$

Combine with (4): $dG = -SdT + Vdp + \mu dN$

gives $N d\mu = -SdT + Vdp$

so Gibbs-Duhem:

$d\mu = -sdT + vdp$ with $s = \frac{S}{N}$
 $v = \frac{V}{N}$

b) $N = -\left(\frac{\partial S}{\partial \mu}\right)_{V, T, A}$

$= V\left(\frac{\partial p}{\partial \mu}\right)_T - A\left(\frac{\partial \gamma}{\partial \mu}\right)_T$

$= A \int dz \rho(z) = \rho_b V + A \int dz (\rho(z) - \rho_b)$

Since $\left(\frac{\partial p}{\partial T}\right)_T = \rho_b$ we have

$\left(\frac{\partial \gamma}{\partial \mu}\right)_T = - \int_0^{\infty} dz (\rho(z) - \rho_b)$

2) 1) $g(z) = \rho_+(z) - 2\rho_-(z)$

$\rho_+(z) = \rho_s e^{-\beta e \psi(z)}$; $\rho_-(z) = \left(\frac{\rho_s}{2}\right) e^{+2\beta e \psi(z)}$

Poisson: $\beta e \psi''(z) = -4\pi \lambda_D^{-2} \left[e^{-\beta e \psi(z)} - 2 \cdot \frac{1}{2} e^{+2\beta e \psi(z)} \right]$ PB

$$\text{BC's: } \psi(\infty) = 0 \quad | \\ \psi(0) = \psi_0 \quad |$$

b) $\psi_0 \ll 10 \text{ mV}$ so $\beta e \psi_0 \ll \frac{10}{25} < 1$ so linearisation is OK. |

$$\beta e \psi''(z) = -n\pi \lambda_B \epsilon_s [1 - \beta e \psi(z) - (1 + 2\beta e \psi(z)) \psi(z)]$$

$$= \beta e n\pi \lambda_B \cdot 3\epsilon_s \psi(z) \quad |$$

$$\Rightarrow \psi''(z) = k^2 \psi(z) \text{ with } k^2 = 12\pi \lambda_B \epsilon_s \quad |$$

$$\text{so } k^{-1} = \frac{1}{\sqrt{12\pi \lambda_B \epsilon_s}} \quad |$$

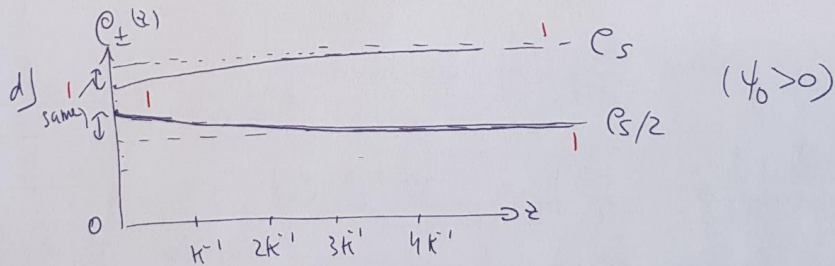
$$\text{c) } \psi(z) = A e^{+kz} + B e^{-kz} = \psi_0 e^{-kz} \quad |$$

$\begin{array}{c} \parallel \\ 0 \\ \text{(BC at } \infty) \end{array}$
 $\begin{array}{c} \parallel \\ \psi_0 \\ \text{(BC at } z=0) \end{array}$

$$\text{Gauss: } \beta e \psi'(0^+) = -n\pi \lambda_B \sigma \quad |$$

$$\parallel \\ -k\psi_0$$

$$\Rightarrow e\sigma = e \frac{k\psi_0}{4\pi \lambda_B} \quad |$$



The positive surface charge is screened by the net negative cloud of ions, of thickness $\sim k^{-1}$ |

$$\text{e) } \lambda_B = 0.7 \text{ nm} \quad | \quad (\text{anything } \sim 1 \text{ nm is fine})$$

iv) $k^{-1} \approx 10 \text{ nm}$ of $\rho_s \approx 1 \text{ m}^{-1}$ (actually a factor $\sqrt{\frac{2}{\pi}}$ smaller)

2

$$k^{-1} \approx 8 \text{ nm}$$

$$\begin{aligned} \text{(ii)} \quad \rho_s \times \text{nm}^3 &= 10^{-3} \cdot (0.6) \cdot 10^{24} \frac{\text{nm}^3}{10^{-3} \text{m}^3} \\ &= (0.6) \cdot 10^{-3} \end{aligned}$$

so volume per ion in nm^3

$$V/\text{nm}^3 = \frac{1}{\rho_s \text{ nm}^3} = \frac{10^3}{0.6} \quad 2$$

3) a) $\phi(r) \equiv 0$: ideal gas

$$Z = \frac{V^N}{N! \Lambda^{3N}} \quad \text{with } \Lambda = \sqrt{\frac{h^2}{2\pi m k_B T}}$$

$$\text{(or } Z = \frac{1}{N! h^{3N}} \int d\mathbf{r}^N \int d\mathbf{p}^N e^{-\beta K} = \dots)$$

$$U = \langle H \rangle = \langle K \rangle = \frac{3}{2} N k_B T \quad 2$$

↑
equipartition

$$P = - \frac{\partial}{\partial V} (-k_B T \ln Z) = \frac{N k_B T}{V} \quad 1$$

$$S = \frac{-F + U}{T} = \frac{k_B T \ln Z + U}{T} \quad \text{with } Z \text{ and } U \text{ as above}$$

$$\mu = \partial F / \partial N = k_B T \ln(N \Lambda^3 / V) \quad 1$$

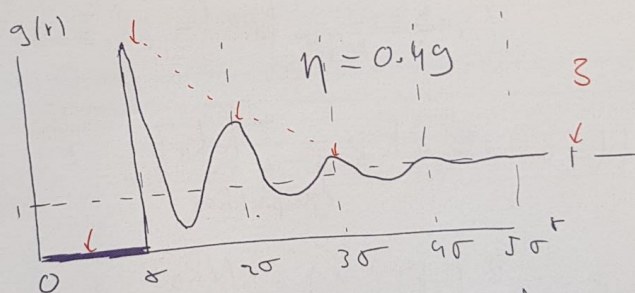
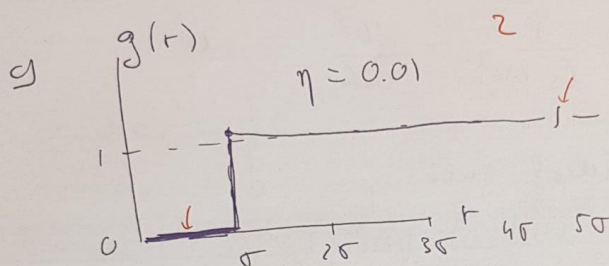
(b) fluid is isotropic & homogeneous

so $\rho^{(2)}(\vec{r}, \vec{r}')$ only depends on $|\vec{r} - \vec{r}'|$

$$\langle \Phi \rangle = \left\langle \sum_{i < j} \phi(r_{ij}) \right\rangle =$$

$$= \frac{1}{2} \left\langle \sum_{i=1}^N \sum_{j=1}^N \left(d\mathbf{r} d\mathbf{r}' \delta(\vec{r} - \vec{r}') \delta(\mathbf{r}' - \vec{r}_j) \phi(|\vec{r} - \vec{r}'|) \right) \right\rangle$$

$$\begin{aligned}
 &= \frac{1}{2} \left\langle \sum_{i=1}^N \sum_{j \neq i}^N \int d\vec{r} d\vec{r}' \delta(\vec{r} - \vec{r}') \delta(r' - r_j) \phi(|\vec{r} - \vec{r}'|) \right\rangle \\
 &= \frac{1}{2} \int d\vec{r} d\vec{r}' \rho^{(2)}(\vec{r}, \vec{r}') \phi(|\vec{r} - \vec{r}'|) \\
 &= \frac{1}{2} \rho^2 V \int d\vec{r}_2 g(r_2) \phi(r_2) \\
 &= \frac{1}{2} \rho^2 V 4\pi \int_0^\infty dr r^2 g(r) \phi(r)
 \end{aligned}$$



In both cases $\langle \Phi \rangle = 0$ because
 either $g(r) = 0$ ($r < \sigma$)
 or $\phi(r) = 0$ ($r > \sigma$)

$$\begin{aligned}
 d) \quad B_2 &= \frac{1}{2} \int d\vec{r} (1 - e^{-\beta\phi(r)}) \\
 &= \frac{1}{2} 4\pi \left(\int_0^\sigma dr r^2 + \int_\sigma^{2\sigma} dr r^2 (1 - e^{-\beta\phi(r)}) \right)
 \end{aligned}$$

$$= \frac{2\pi}{3} \sigma^3 + \frac{2\pi}{3} (1 - e^{\beta \epsilon}) \sigma^3 (\rho - 1)$$

$$= \frac{2\pi}{3} \sigma^3 + \frac{14}{3} \sigma^3 (1 - e^{\beta \epsilon}) \quad 2$$

$$e) H_\lambda = H_0 + \lambda \Phi_1, \quad H_0 = K + \sum_{i,j} \phi_0, \quad \Phi_1 = \sum_{i,j} \phi_1(r_{ij})$$

$$\frac{\partial F_\lambda}{\partial \lambda} = \langle \Phi_1 \rangle_\lambda$$

$$\phi_1(r) = \begin{cases} -\epsilon & \sigma < r < 2\sigma \\ 0 & \text{otherwise} \end{cases}$$

$$\text{So } F_{\lambda=1} \equiv F$$

$$F_{\lambda=0} + \int_0^1 d\lambda \langle \Phi_1 \rangle_\lambda =$$

$$= F_{HS} + \int_0^1 d\lambda \frac{1}{2} \rho^2 V \int_0^\infty dr 4\pi r^2 \phi_1(r) g_\lambda(r)$$

$$= F_{HS} + \frac{1}{2} \rho^2 V 4\pi \int_\sigma^{2\sigma} dr r^2 (-\epsilon) \int_0^1 d\lambda g_\lambda(r)$$

f) 1st order perturbation theory: $g_\lambda(r) \approx g_0(r)$ 2

$$\text{So } a = 2\pi \epsilon \int_\sigma^{2\sigma} dr r^2 g_0(r) \quad 1$$

$\rho \sigma^3 = 0.4$: no crystallisation but

at low enough $T < T_c$ VdW loop develops in free energy \rightarrow Gas-Liquid transition 1

\rightarrow system is not stable at all T . 1

4) ... $\nu = 1/2 > D$ so R_n is swollen

$$g = \frac{bN}{2} \sim \lambda g$$

Compared to ideal chain.

(ii) monomer volume fraction inside single chain is

$$\phi^* = \frac{Nb^3}{R_g^3} = \frac{N}{N^{3 \cdot 3/5}} = N^{-4/5}$$

b) Virial expansion:

$$\beta \Pi b^3 = b^3 \left[\frac{M}{V} + B_2 \left(\frac{M}{V} \right)^2 + \dots \right]$$

$$\text{with } B_2 = c R_g^3$$

$$\frac{Mb^3}{V} = \phi/N$$

$$\begin{aligned} \text{so } \beta \Pi b^3 &= \phi/N + c R_g^3 b^3 (\phi/N)^2 \\ &= (\phi/N) (1 + c (\phi/\phi^*) + \dots) \end{aligned}$$

$$\text{so } f(x) = 1 + cx \quad \text{if } x \ll 1$$

Defines: $x \gg 1$ then Π should be independent of N ,

$$\text{so } (\phi/N) (N^{4/5} \phi)^m \text{ ind of } N$$

$$\Rightarrow m = 5/4$$

$$\text{and } \Pi \sim \phi^{9/4}$$

$$c) \langle X^2 \rangle = 2D t \quad \text{with } D = \frac{k_B T}{6\pi\eta a} \quad (\text{SE})$$

$$\text{so } t_D = \frac{a^2}{2D} = \frac{12\pi\eta a^3}{k_B T} \sim a^3$$

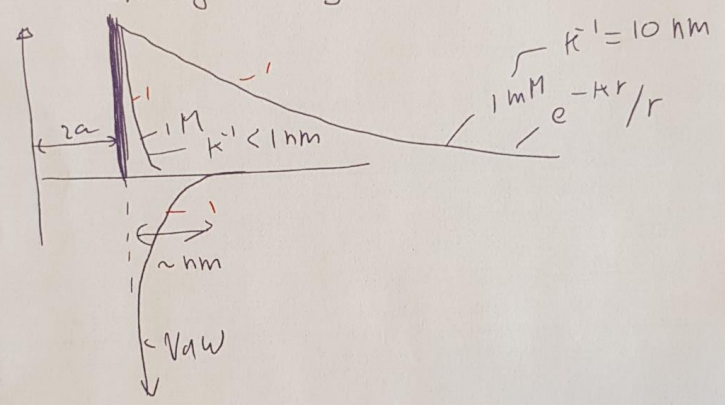
exp techniques: - confocal microscopy

Tentam
Opleidi
Naam:

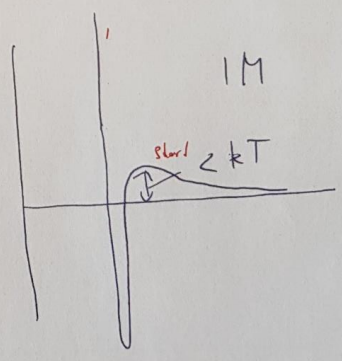
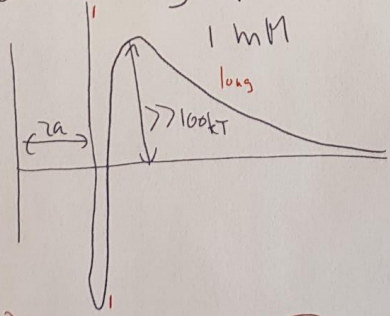
- for $a = 1 \mu\text{m}$
- light scattering
 - (video tracking)
 - - -

d) 3 ingr edients of DLVO :

- steric "hard-core" repulsions
- VDW "dispersion" attractions
- Screened-Coulomb repulsions due to charged surfaces & ionic screening



So summing up



total 3 points